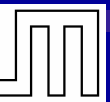


THE SIMULATION OF MOLECULAR AND ORGANIC DEVICES: A CRITICAL REVIEW AND A LOOK AT FUTURE DEVELOPMENTS

P. Lugli

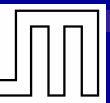
Lehrstuhl für Nanoelektronik
Technische Universität München

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A. Bolognesi, A. Di Carlo, A. Pecchia, L. Latessa (Rome-“Tor Vergata”)

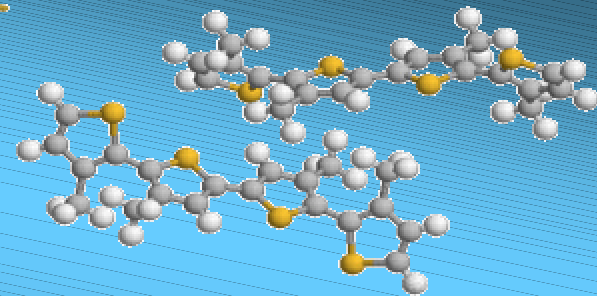
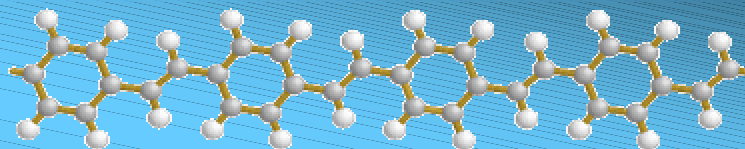
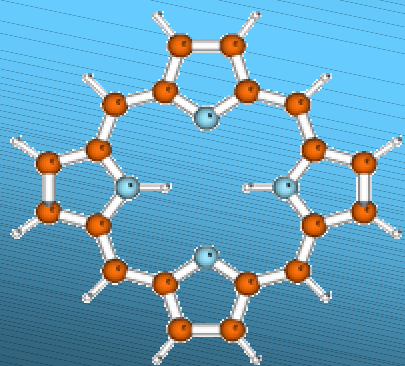


Outline

- Overview of molecular electronics
 - Generalities
 - OTFT Simulation
- Single molecule electronics
 - DFTB method
 - Simulation of molecular diodes
- Carbon Nanotubes
 - CNTs as sensors and transistors
- Open questions and closing remarks



Approaches to molecular electronics



Organic Molecules
100-1000 atoms or more

Organic Material for electronics

Conventional simulators

- Organic Transistors
- Organic LED
- Organic Solar Cells

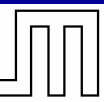
Conventional devices with organic semiconductors

Molecular Electronics

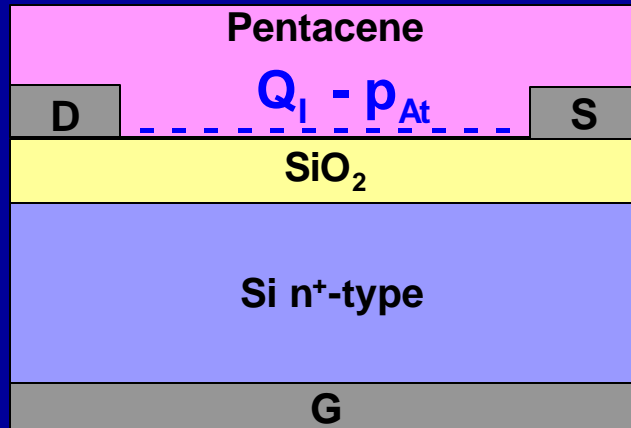
New simulators

- Diodes
- Transistors
- Memories

New devices with nanosize dimension



Simulation Model



The simulation uses a modified drift-diffusion tool (ISETCAD).

Trap model according to Scheinert et al. JAP 92, 330 (2002)

$$m = m_0(T) \exp\left(\sqrt{\frac{E}{E_0}}\right)$$

$$\nabla_n = -q\mathbf{m}_n \nabla V + qD_n \nabla n$$

$$\nabla_p = -q\mathbf{m}_p \nabla V - qD_p \nabla p$$

$$\nabla \cdot (\mathbf{e}_0 \mathbf{e}_r \nabla V) = -q(p - n + N_D^+ - N_A^- - p_{At}) - Q_I$$

$$\nabla \cdot \mathbf{J}_n = q(R - G)$$

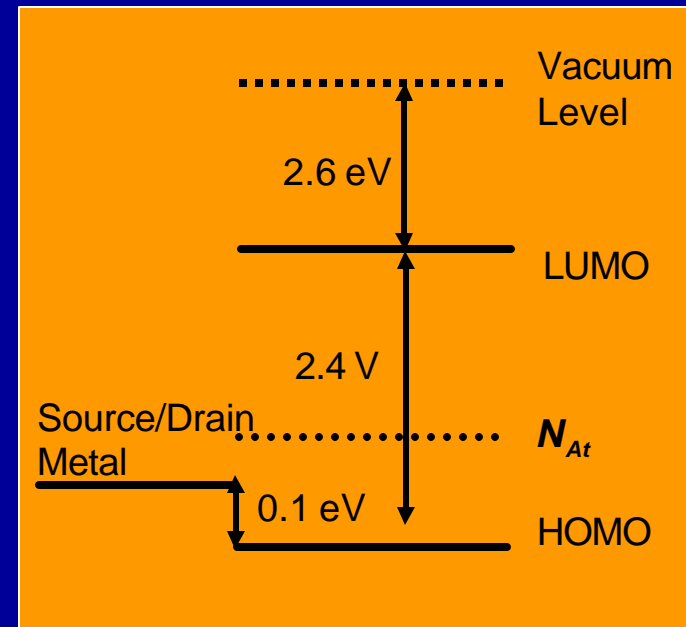
$$\nabla \cdot \mathbf{J}_p = -q(R - G)$$

$$p_{At} = N_{At} f_n$$

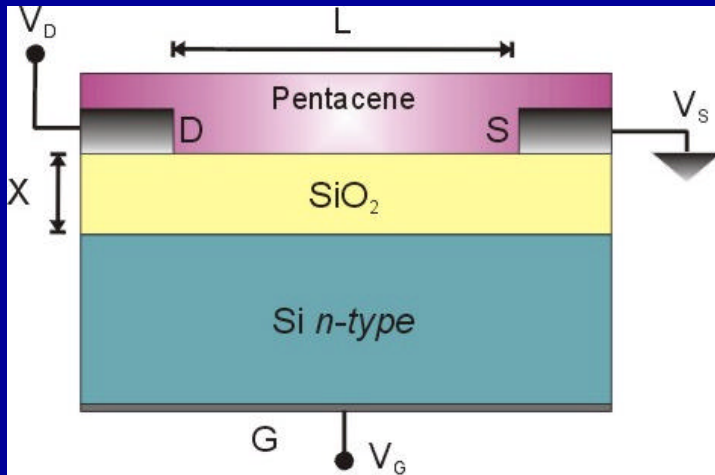
Acceptor-like traps

$$Q_I \quad \text{SiO}_2/\text{Pentacene interface charge}$$

Pentacene levels alignment

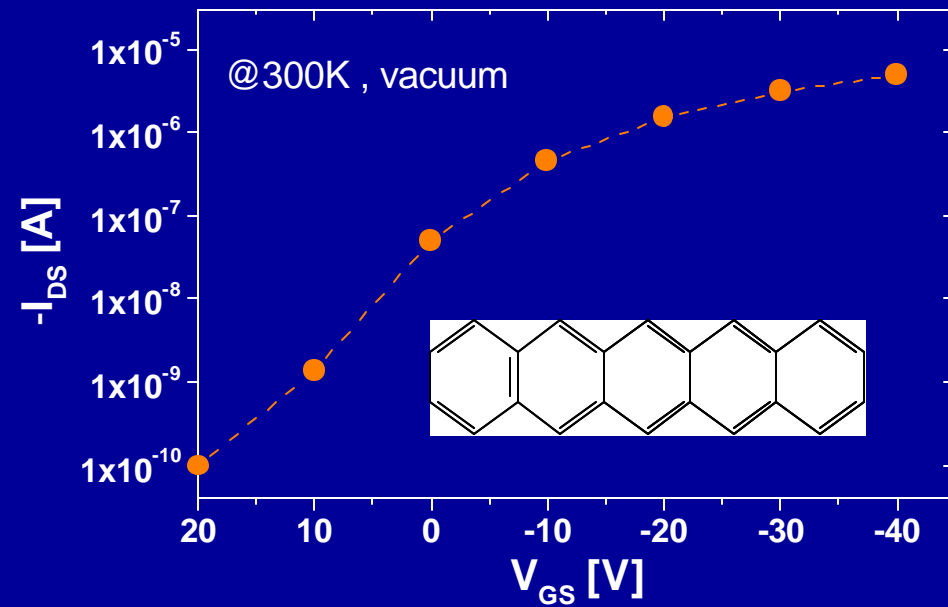
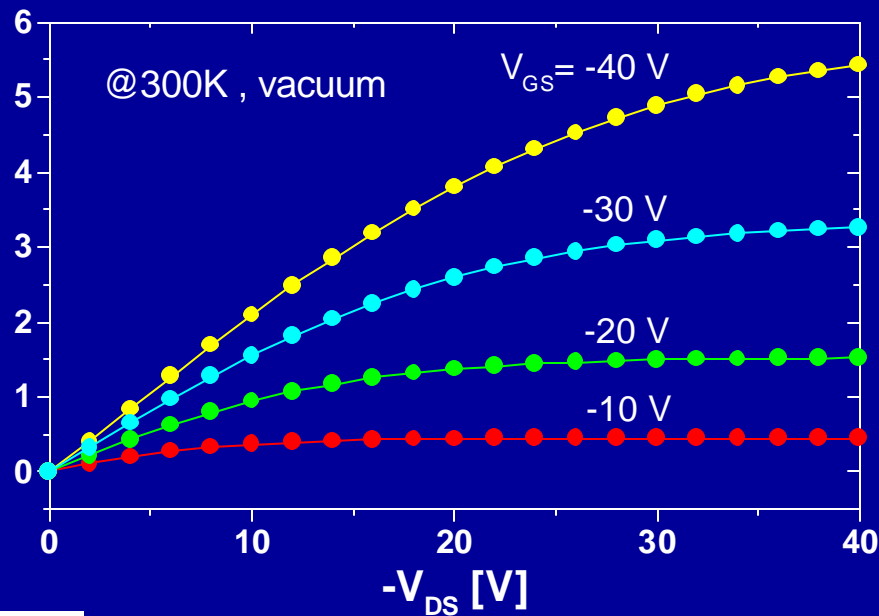


OTFT characteristics

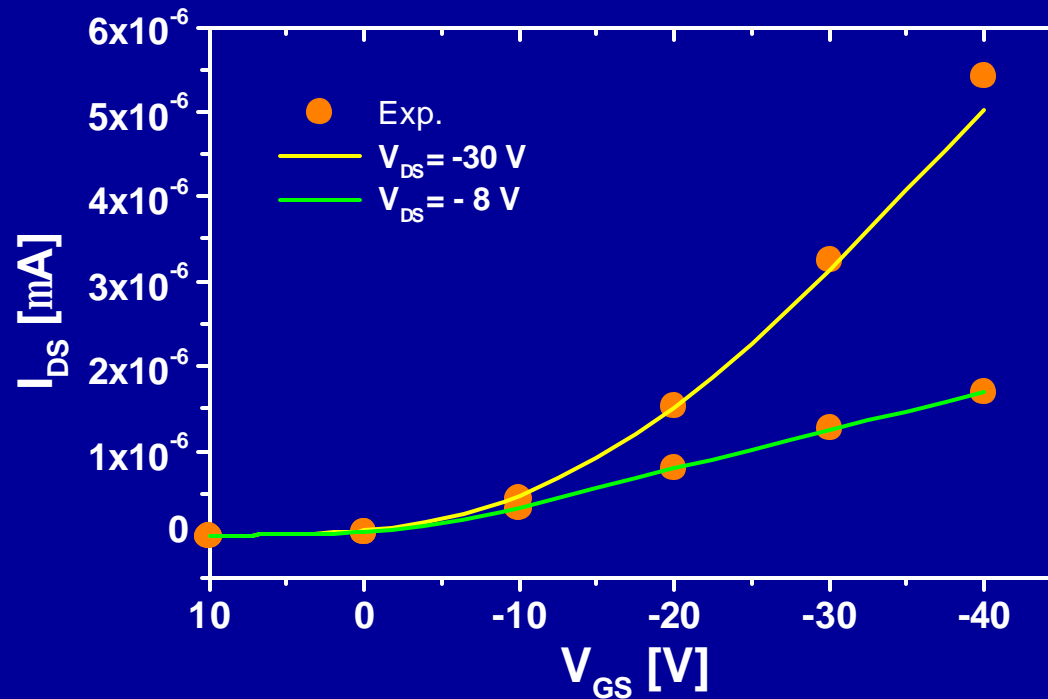


- Thin film transistor were fabricated using bottom contact configuration.
- Cr/Au Source/Drain Contact prepared by vacuum evaporation.
- Pentacene sublimated in vacuum at a rate of 0.2 Å/sec.

Channel Length $L=12\mu\text{m}$ ($L / W = 140$)
Oxide Thickness $x=250\text{ nm}$



Transfer Characteristics

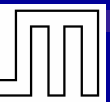


The overall behavior of the device has been simulated and good agreement is found between experiment and simulations.

Parameters: $\mu_0 = 3.3 \text{ cm}^2/\text{Vs}$, $N_I = 2 \cdot 10^{11} \text{ cm}^{-2}$, $N_{\text{trap}} = 10^{12} \text{ cm}^{-2}$, $E_{\text{trap}} = 0.15 \text{ eV}$

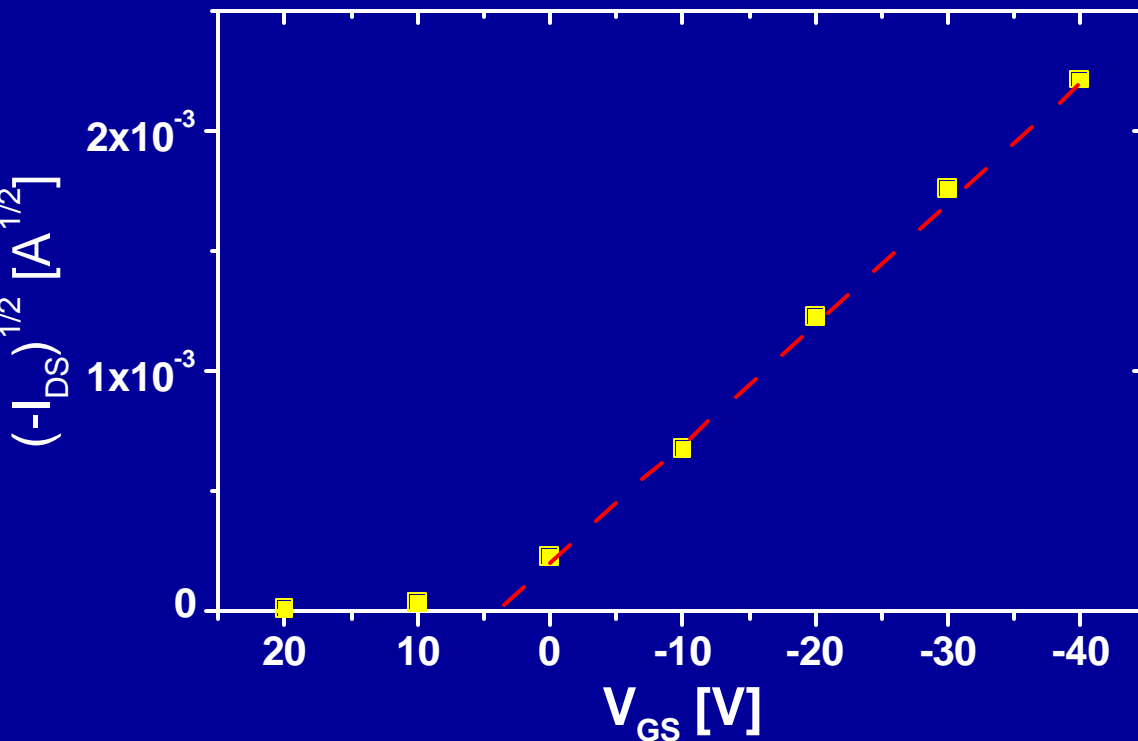
Is this parameterization transferable to other devices ?

YES, if growth conditions remain the same **NO**, if growth conditions are changed



Experimental determination of mobility

Usually standard MOSFET theory is used to extract the mobility from experimental data via the quadratic relation between V_G and I_{Dsat}

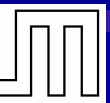


$$I_{Dsat} = m \frac{W}{2L} C_{ox} (V_G - V_T)^2$$

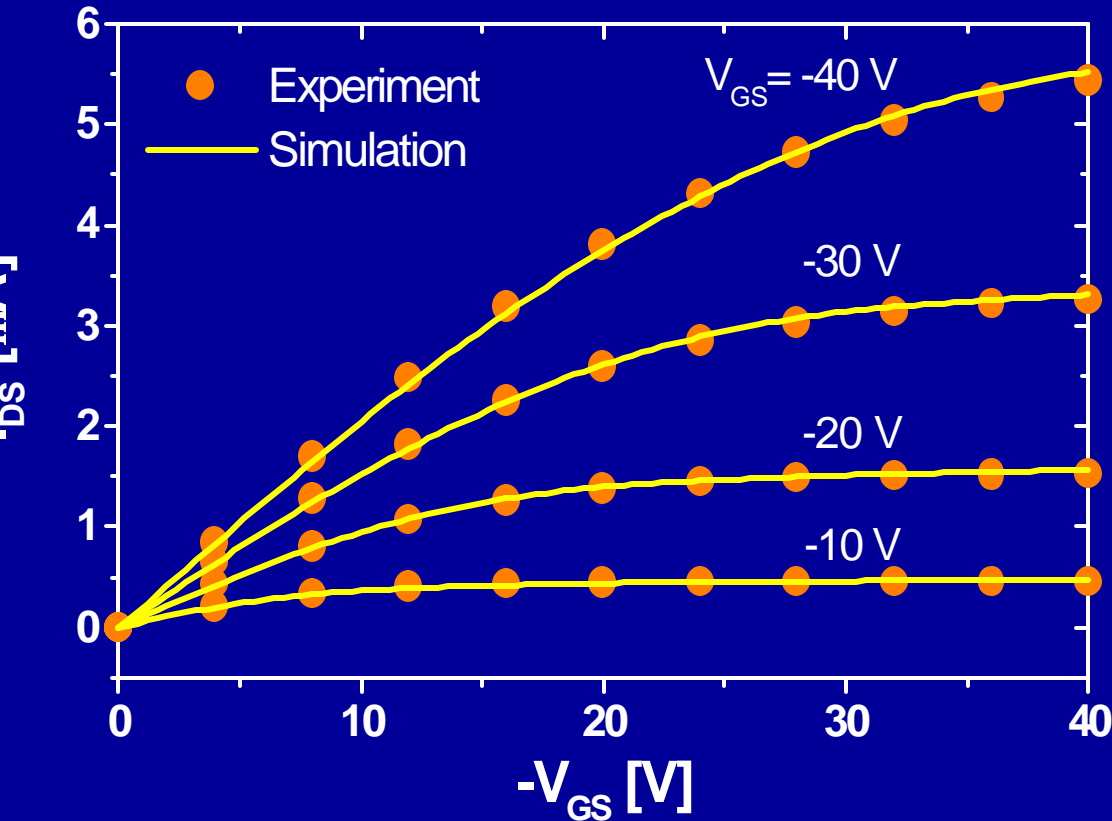


$$m_{exp} = 4.23 \cdot 10^{-3} \frac{\text{cm}^2}{\text{Vs}}$$

How does this mobility compare with the field activated mobility $m = m_0 \exp(\sqrt{\frac{E}{E_0}})$?



Output characteristics: Simulation vs. experiment



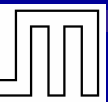
Optimal fitting of the experimental data is found for a

$$m = m_0 \exp\left(\sqrt{\frac{E}{E_0}}\right)$$

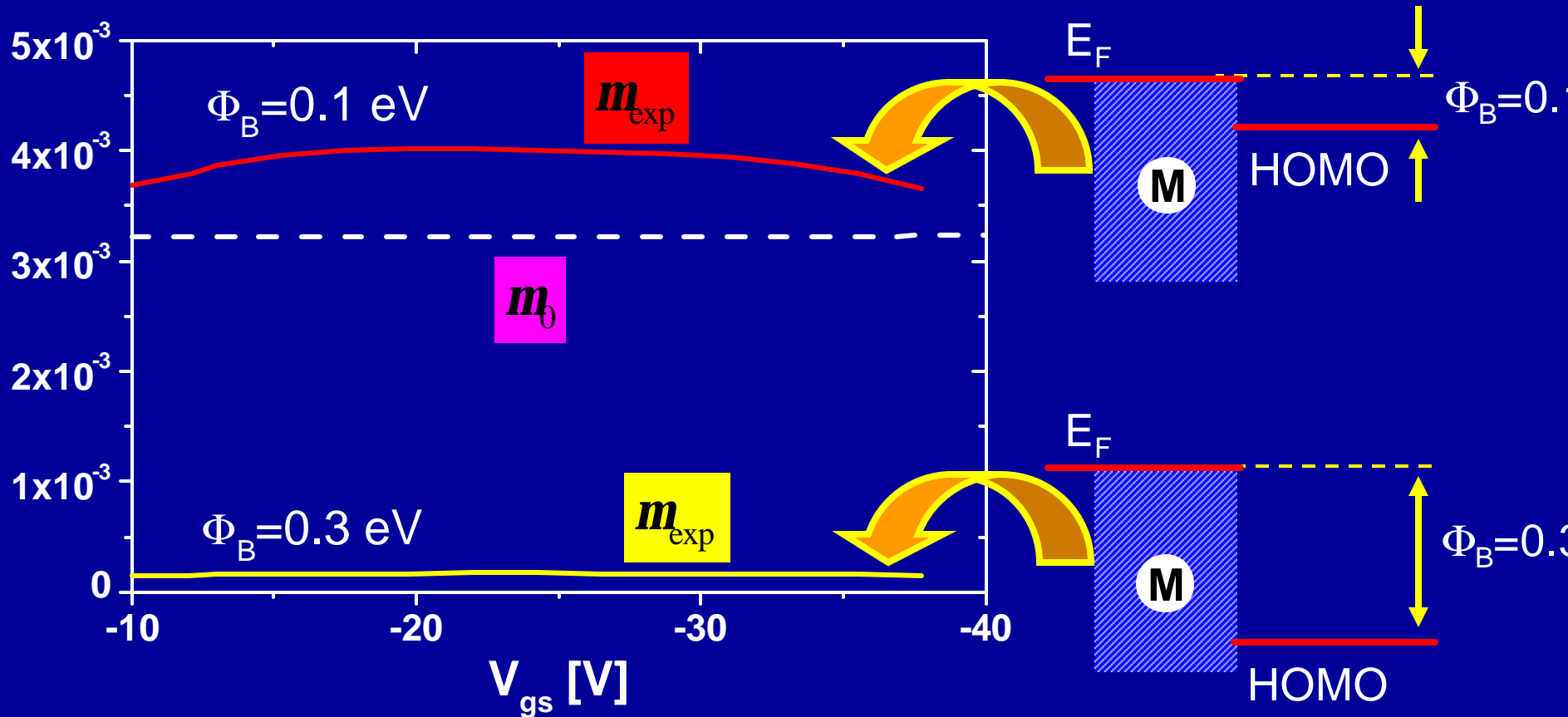
with

$$m_0 = 3.3 \cdot 10^{-3} \frac{\text{cm}^2}{\text{Vs}}$$

$$m_{\text{exp}} \neq m_0$$

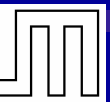


Influence of contact barrier height

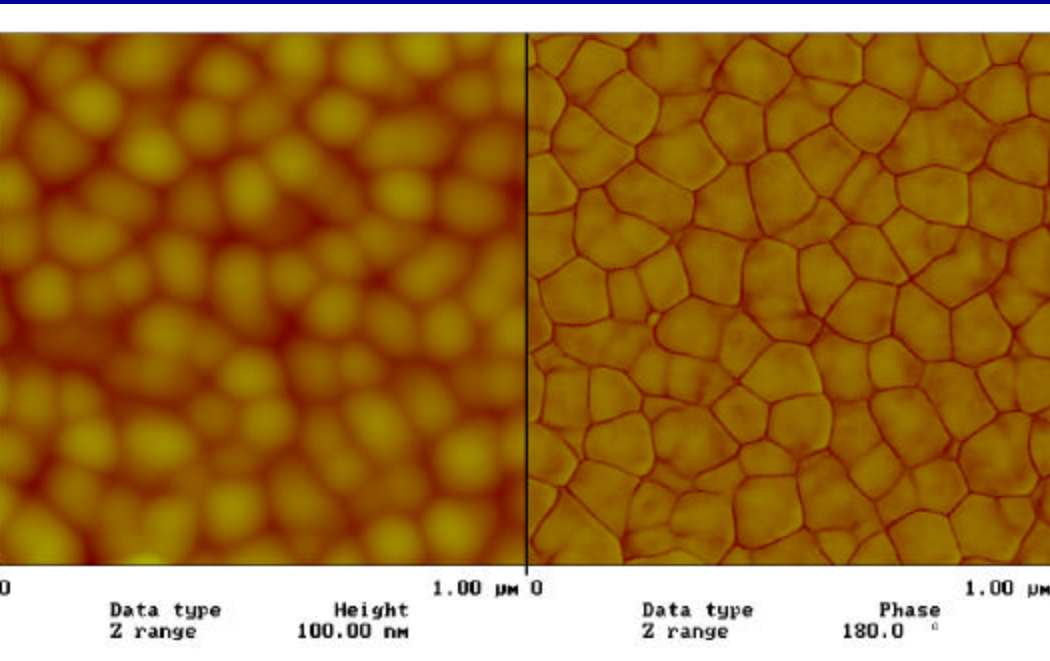


The results are obtained by extracting the mobility from simulation results performed with the field dependent mobility model using different barrier height for the contact

Mobility extraction is strongly influenced by the barrier height.



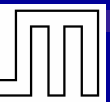
Grain Boundaries



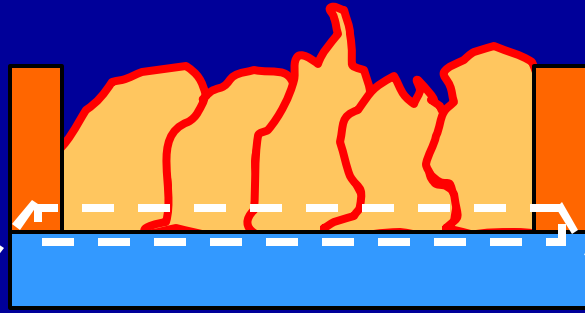
AFM picture of pentacene sample. The size of the grain is about $0.1 \mu\text{m}$.

The concentration of charges trapped at grain boundaries depends on the density of the traps and the position of the Fermi level.

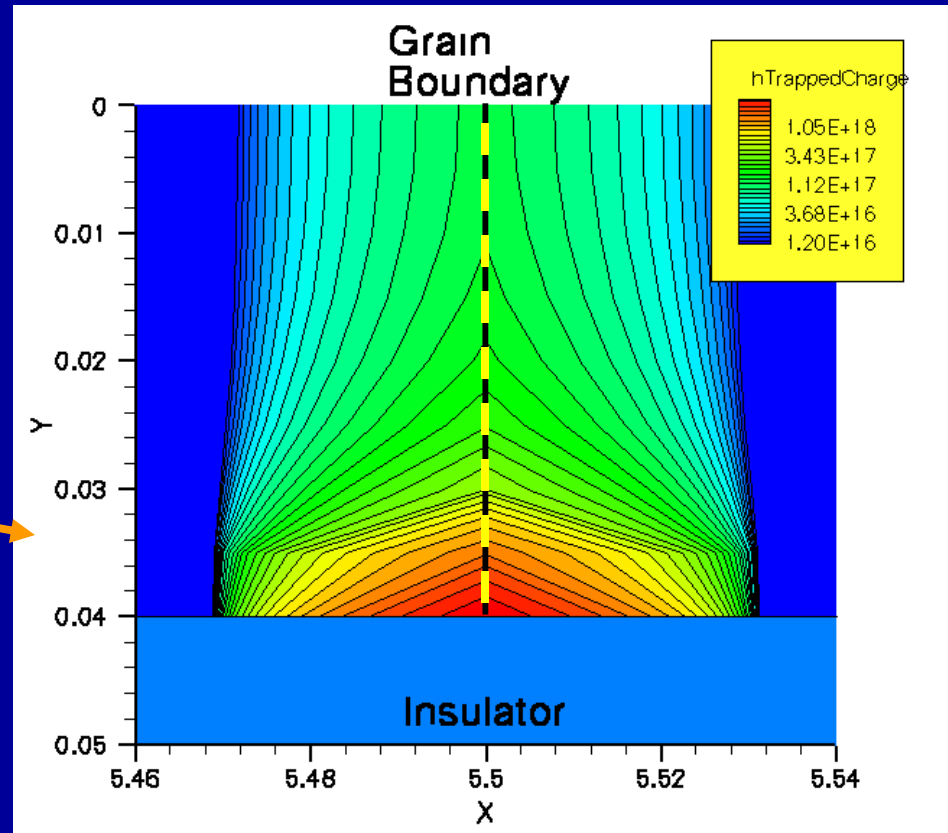
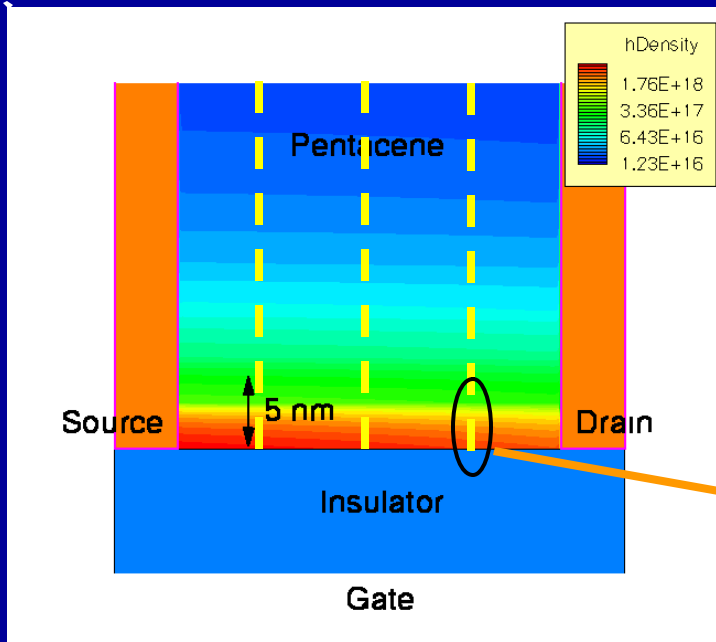
We use single trap level of surface density $N_t=1 \times 10^{13} \text{ cm}^{-2}$ and energy $E_t=0.3 \text{ eV}$



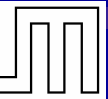
Grain Boundaries



Trapped charge density for $V_{gs} = -20$ V and $V_{ds} = -20$ V

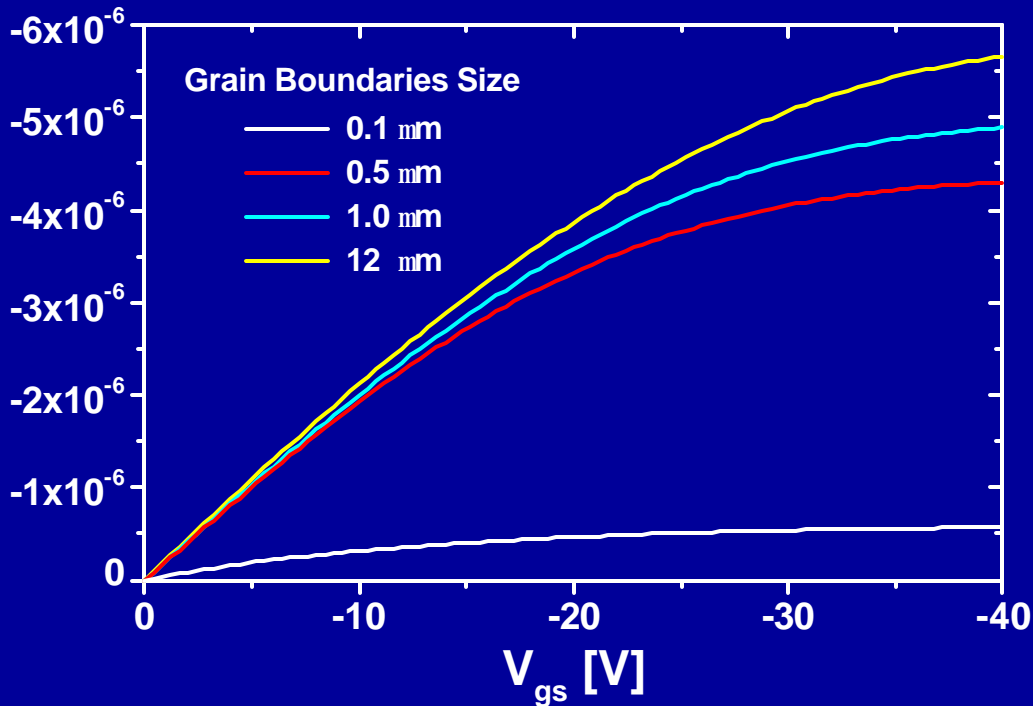


The thickness of conducting channel is lower than 5nm, Thus, it is possible to model the grain with regular shape.

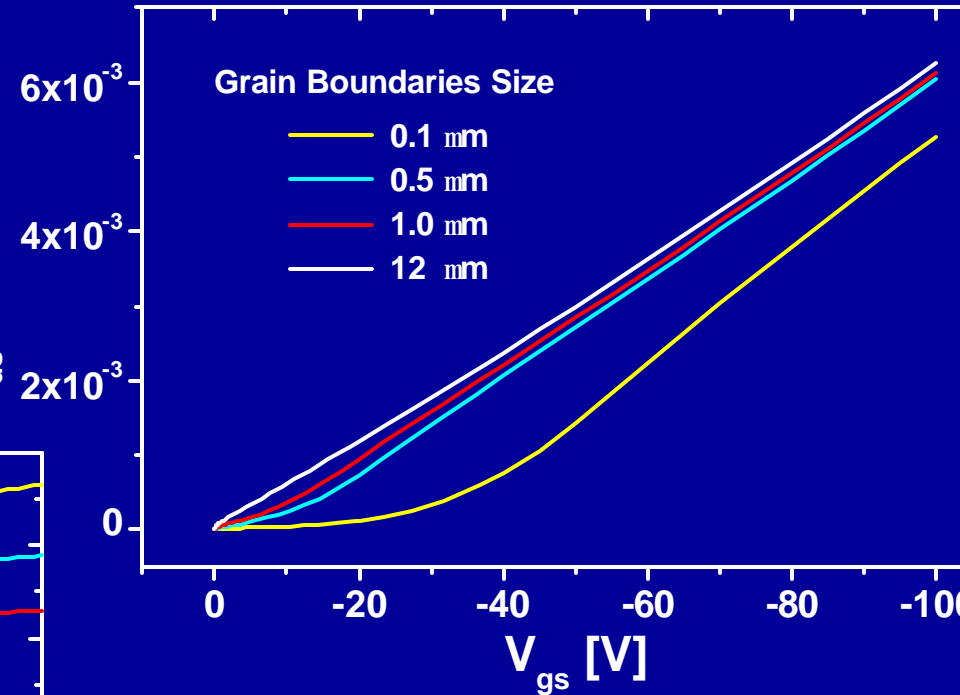


Electrical characteristics

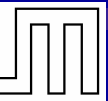
Output characteristics



$I_{ds}^{1/2}$ [A]^{1/2}

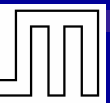


Trans-characteristics



Open problems and perspectives

- ❖ Find a reliable model for carrier mobility and contact injection mechanism
- ❖ Enhance the predictability of the simulation tools
- ❖ Move to physically based circuit simulations



Mobility Model

Field-dependent mobility calculated from MC simulation

- cubic lattice of 170x170x20 hopping sites
- Energies of the sites chosen randomly from a Gaussian DOS
- Under the influence of external field the mean energy is given as

$$E(x) = E_0(x) - qFx$$

- Rate of hopping described by Miller-Abrahams expression

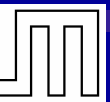
$$G_{i,j} = v_0 \exp(-2gR_{ij}) \exp\left\{-\frac{E_j - E_i + |E_j - E_i|}{2kT}\right\}$$

- MC estimator for the mobility:

$$m = \left\langle \frac{L}{(\dot{a} t_i) F} \right\rangle$$

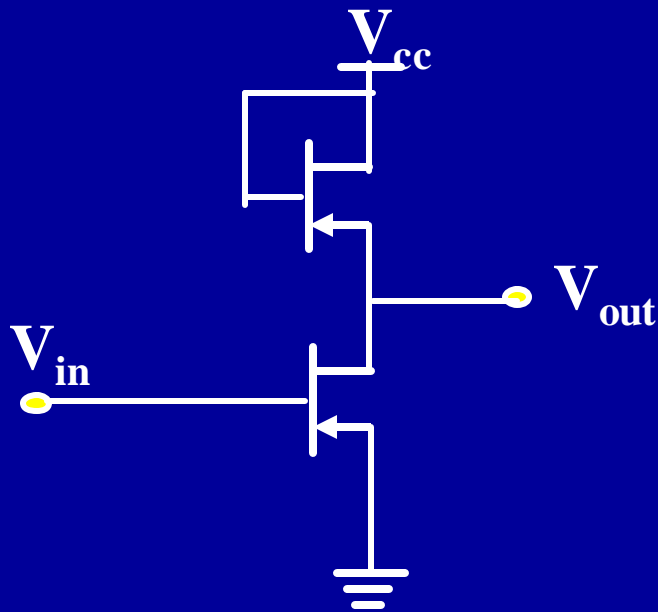
from carrier
lifetime t_i

$$t_i = -\frac{\ln(u)}{\sum_j \Gamma_{ij}}$$



OTFT-based circuits

Inverters and ring oscillators can be simulated directly using the ISETCAD tool

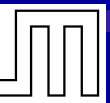
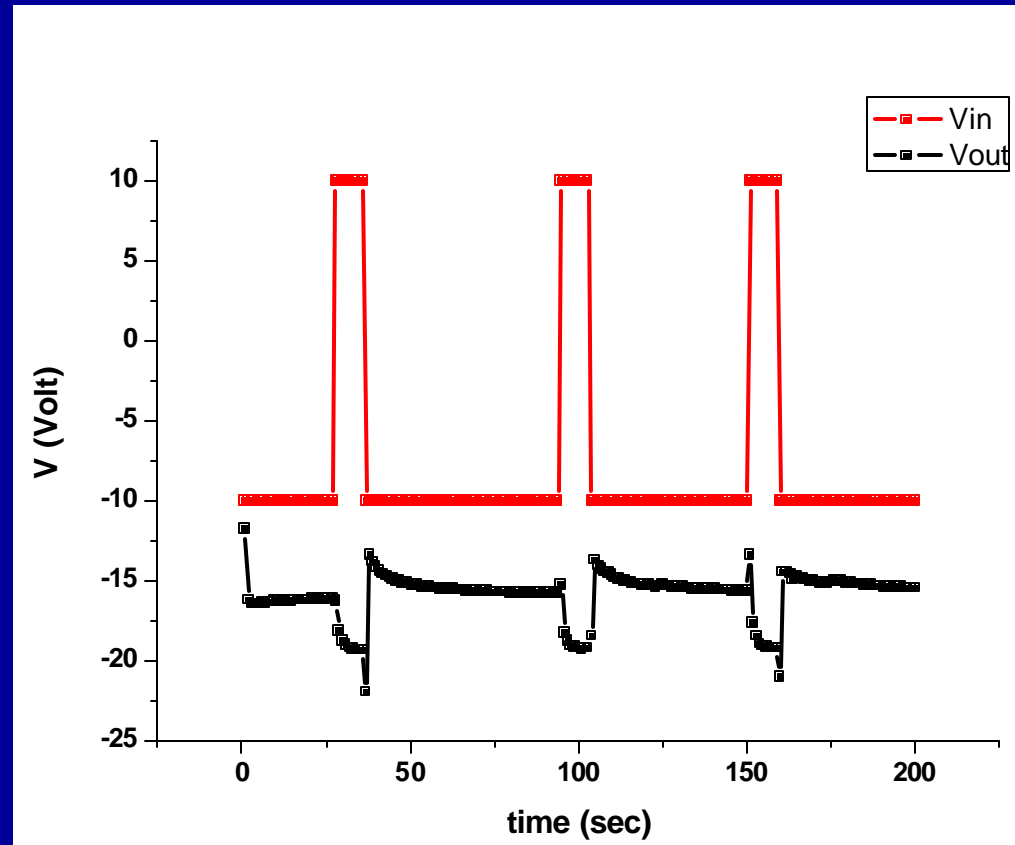


$V_{in_{pp}}: 20\text{ V}$

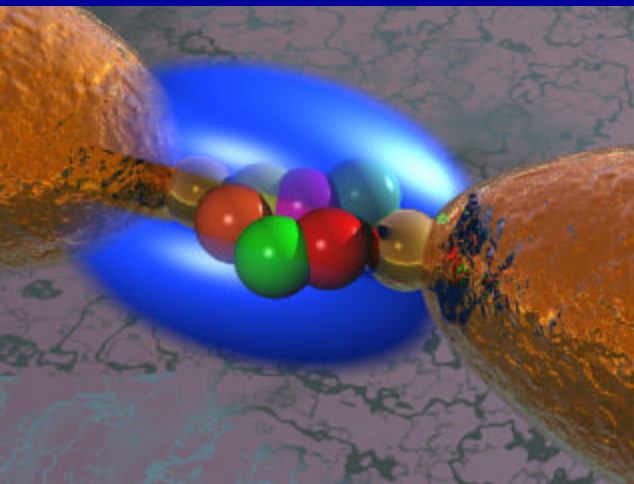
$V_{cc}: -20\text{V}$

Channel length: $20\ \mu\text{m}$

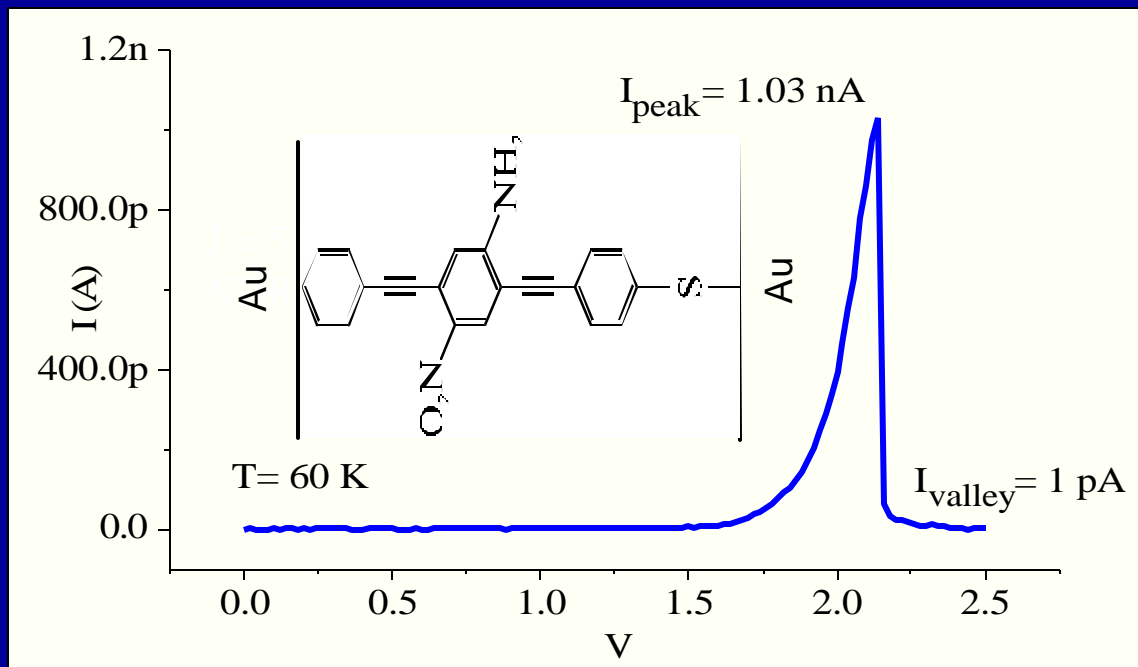
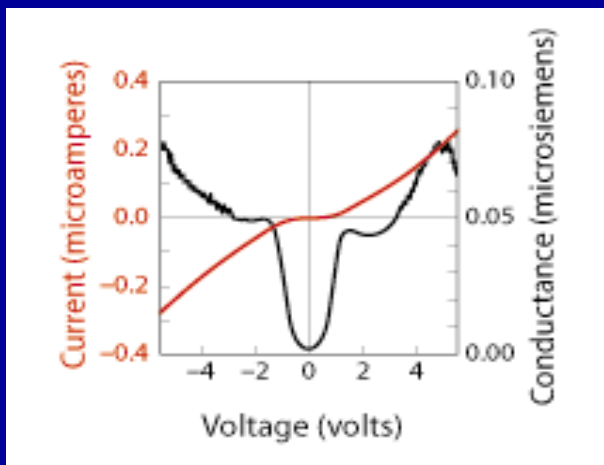
Pentacene



Molecular electronics

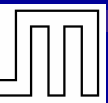


Single molecule can be conductive and display an electrical nonlinear response

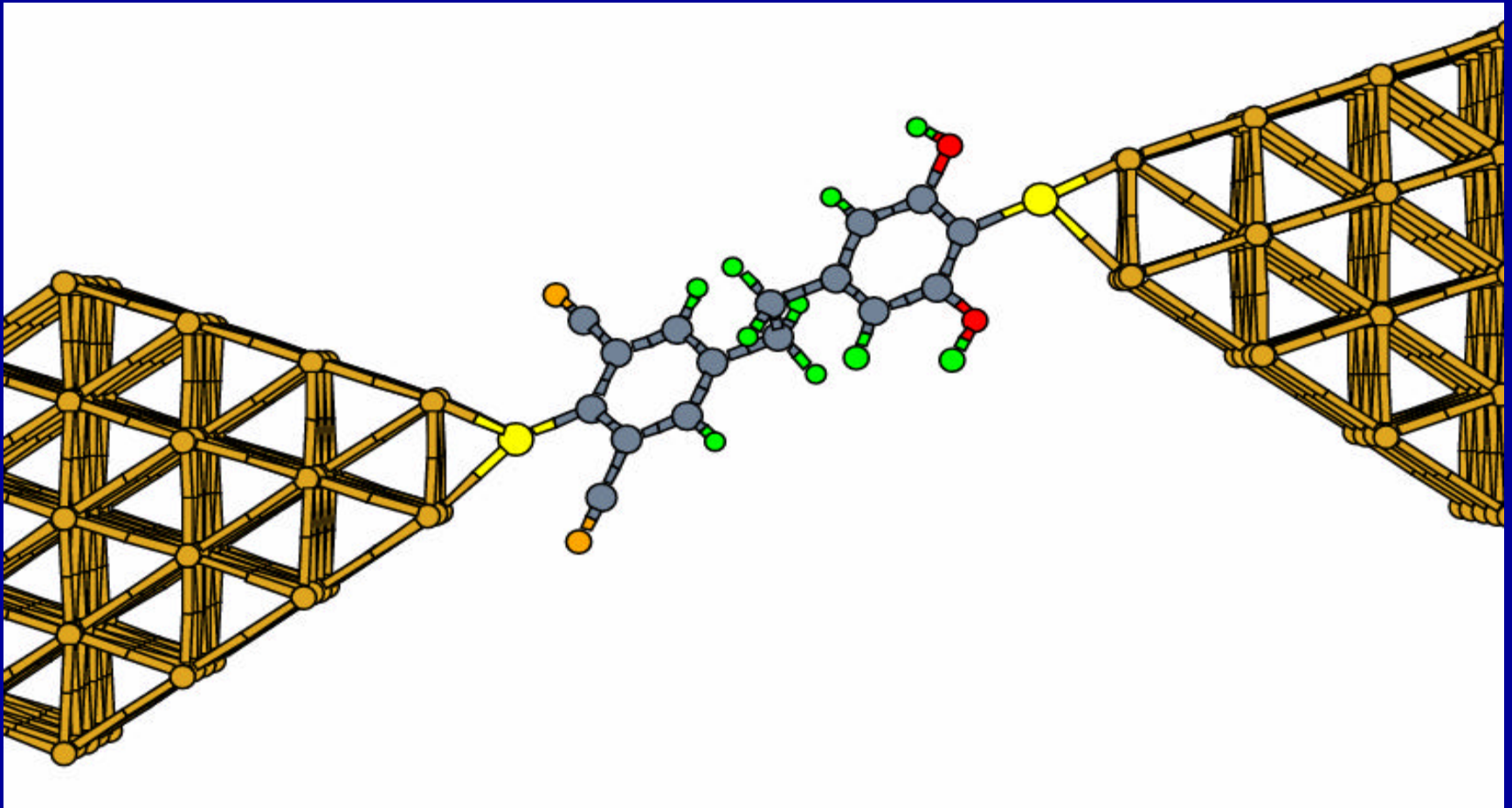


M. Reed et al, Scientific American 2002

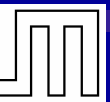
J. Chen et. al, Science 286, 1550-1552 (1999)



Molecular device

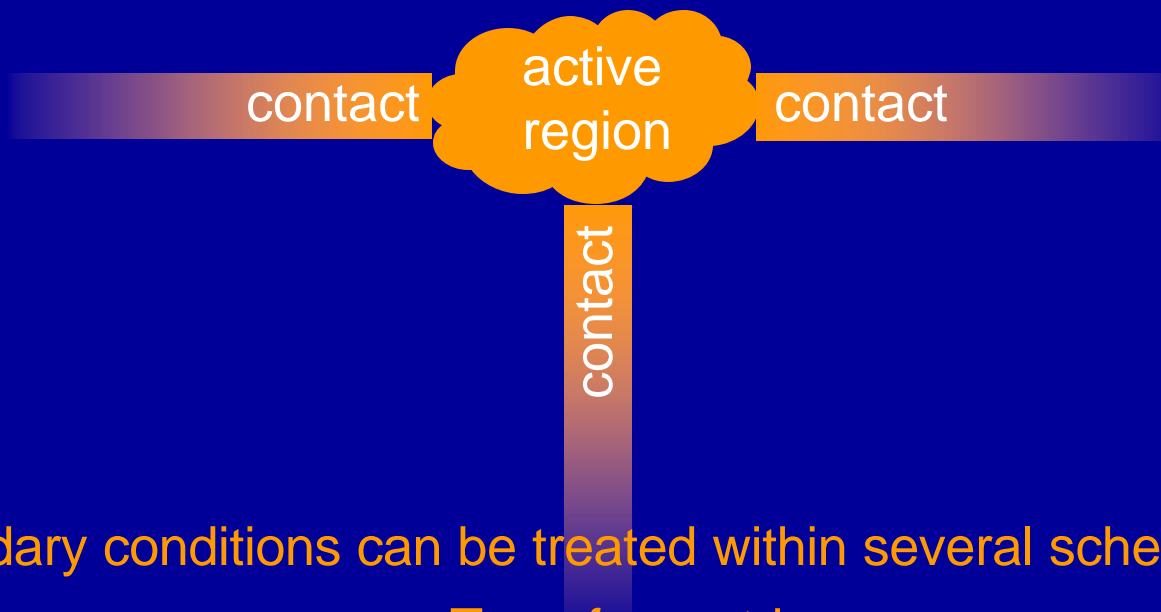


A physically-based method is needed that determines the electronic structure of the molecule and solves the transport problem



Transport problem

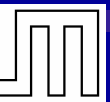
active region where symmetry is lost
+
contact regions (semi-infinite bulk)



Open-boundary conditions can be treated within several schemes:

- Transfer matrix
- LS scattering theory
- Green Functions

These schemes are well suited for localized orbital approach like TB



Atomistic approaches

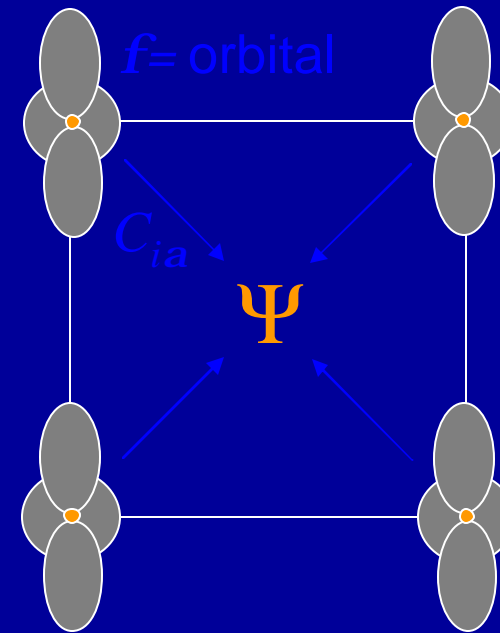
We attempt to solve the one electron Hamiltonian in terms of a Linear Combination of Atomic Orbitals (LCAO)

$$\Psi_n(\mathbf{r}) = \sum_{\substack{\text{atomic} \\ \text{site, } i}} \sum_{\substack{\text{orbitals, } a}} C_{ia} \mathbf{f}_{ia}(\mathbf{r} + \mathbf{R}_i)$$

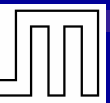
$$\sum_{\substack{\text{atomic} \\ \text{site, } j}} \sum_{\substack{\text{orbitals, } b}} [H_{ia,jb} - ES_{ia,jb}] C_{jb} = 0$$

$$H_{ia,jb} = \langle \mathbf{f}_{ia} | H | \mathbf{f}_{jb} \rangle$$

$$S_{ia,jb} = \langle \mathbf{f}_{ia} | \mathbf{f}_{jb} \rangle$$



The approach can be implemented “ab-initio” where the orbitals are the basis functions and $H_{ia,jb}$ is evaluated numerically



Toward "ab-initio" approaches: Density Functional Tight-Binding

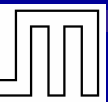
Many DFT codes exist based on localized basis sets:

SIESTA (Soler etc.), **FIREBALL** (Sankey), **DMOL** (Delley)

The **DFTB** approach [Eltner, et al. *Phys. Rev. B* 58 (1998) 7260] provides transferable and accurate interaction potentials. The numerical efficiency of the method allows for molecular dynamics simulations in large super cells, containing several thousand of atoms.

- DFTB is fully scalable (from empirical to DFT)
- DFTB allows also for TD-DFT simulations

We have extended the **DFTB** to account for transport in organic/inorganic nanostructures by using **Non Equilibrium Green Function approach** self-consistently coupled with Poisson equation



Self-consistent quantum transport

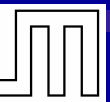
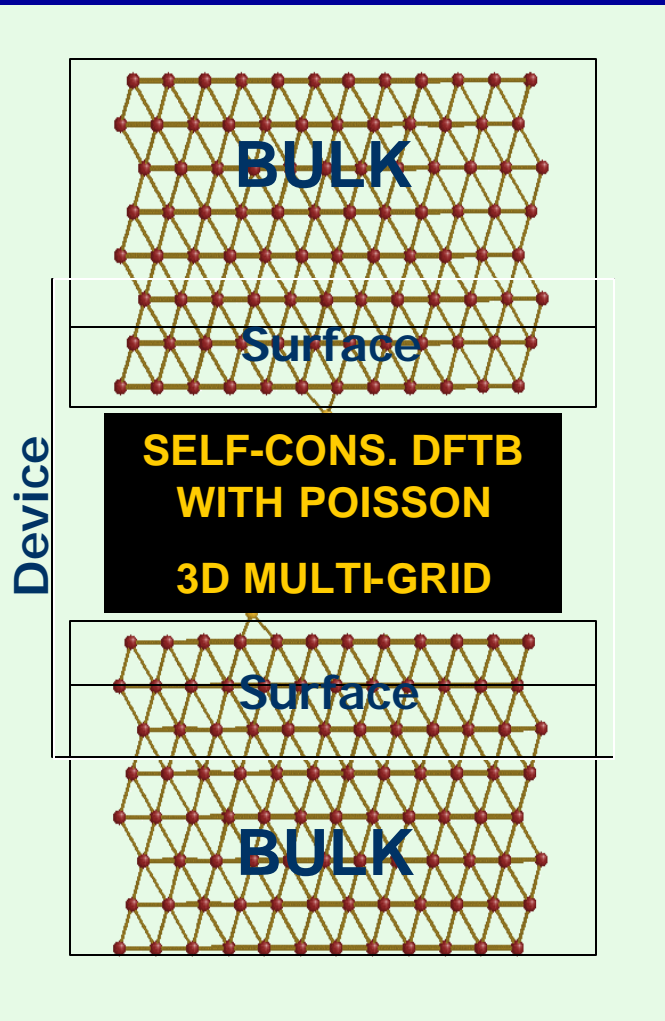
Self-consistent loop:

$$\begin{aligned} d n &\rightarrow \nabla^2 d V_H = -4 p d n \\ &\rightarrow d H \rightarrow d G^n \rightarrow d n' \end{aligned}$$

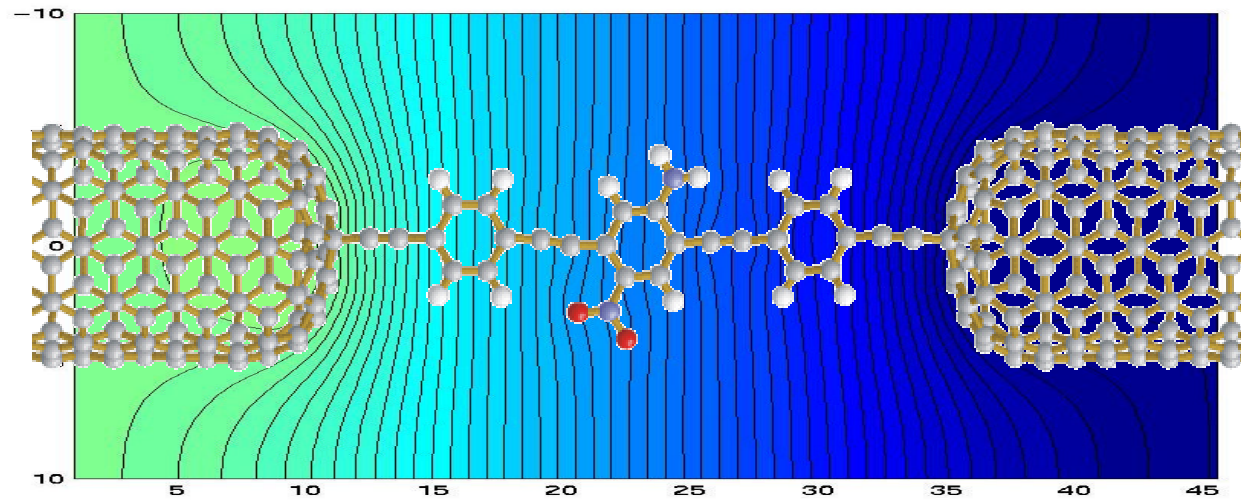
Current computation:

$$T = \text{tr}(\Gamma_L G_D^r \Gamma_R G_D^a)$$

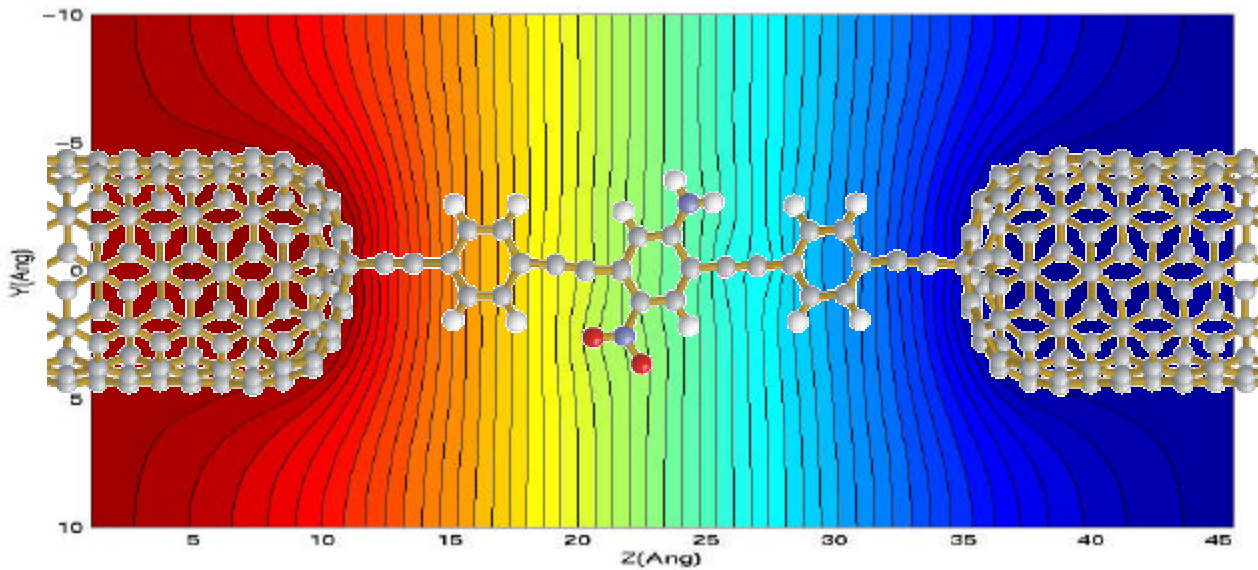
$$i(E) = \frac{2e}{h} \int T(E) [f_L(E) - f_R(E)] dE$$



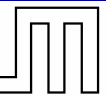
Self-consistent charge in a molecular wire



0.5 V

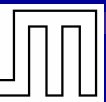
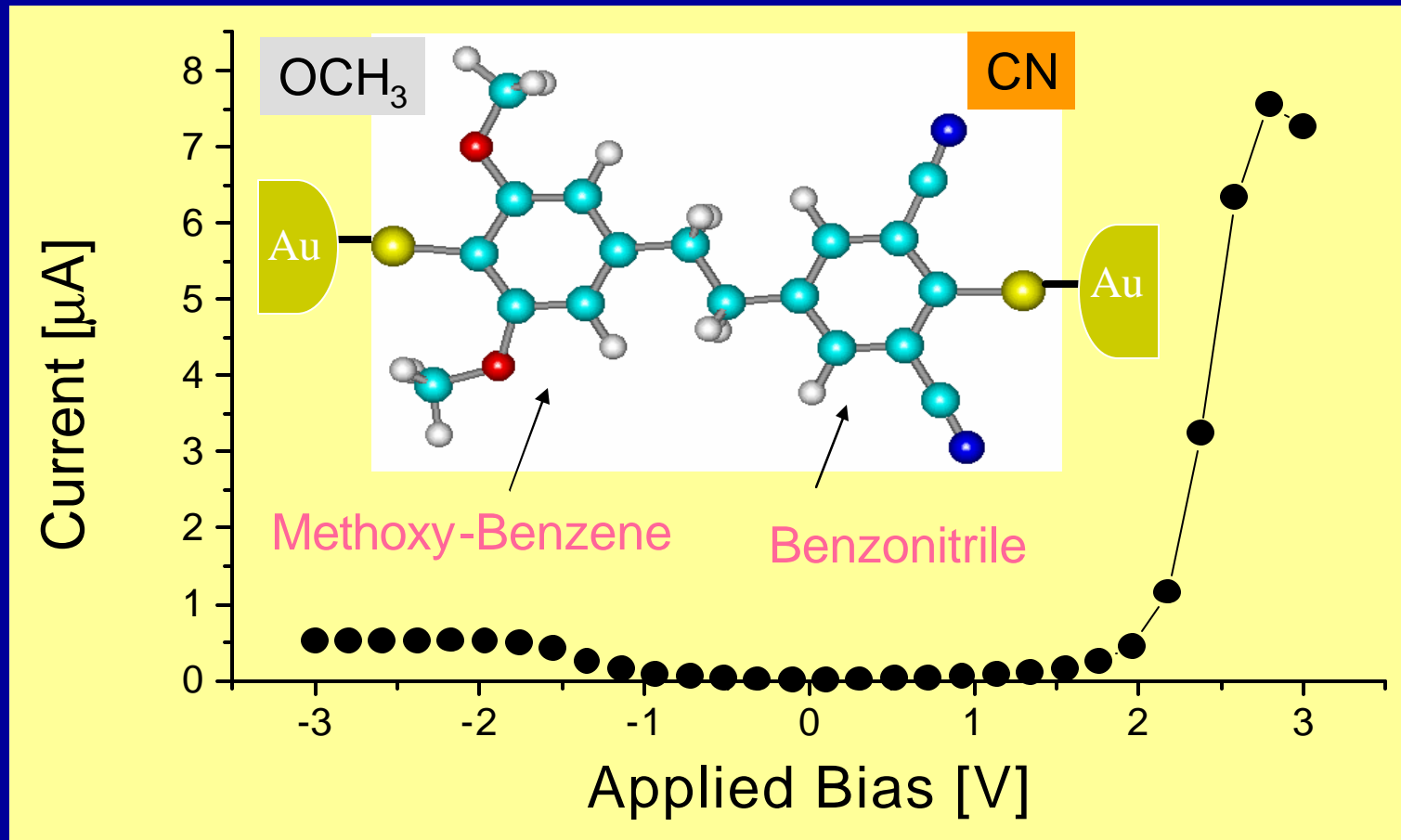


1.0 V



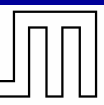
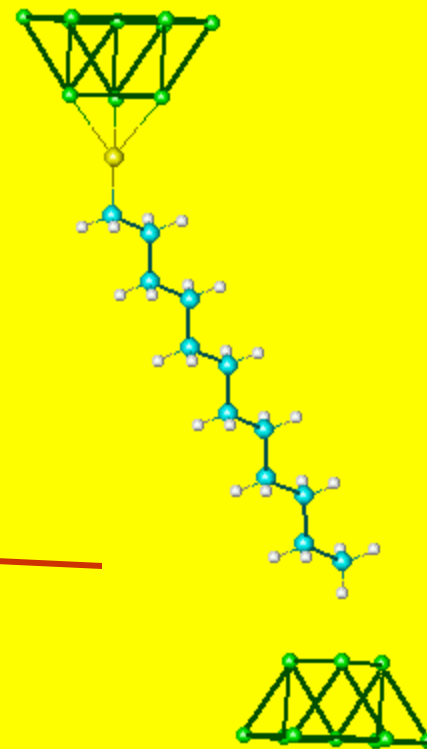
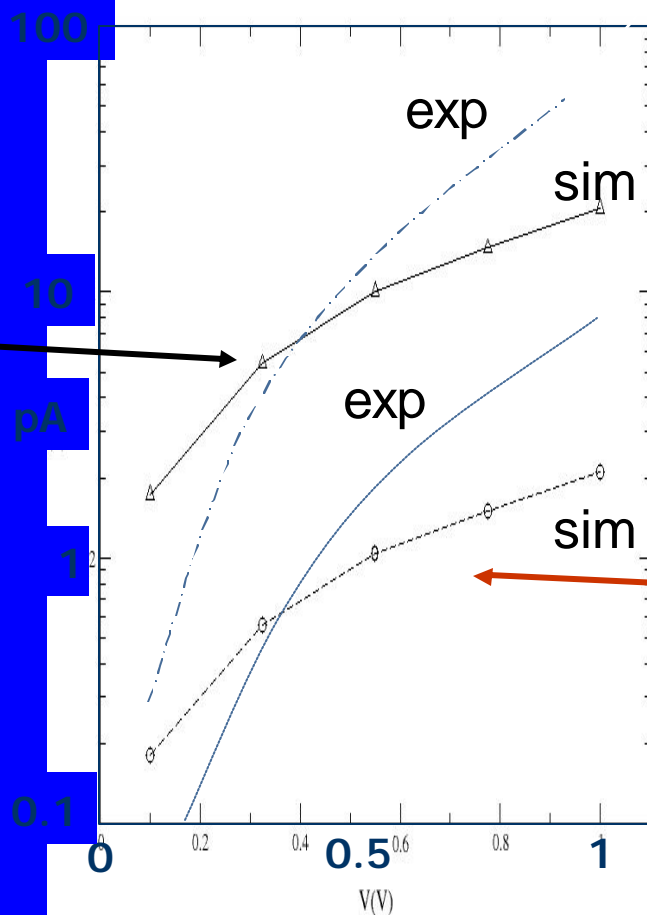
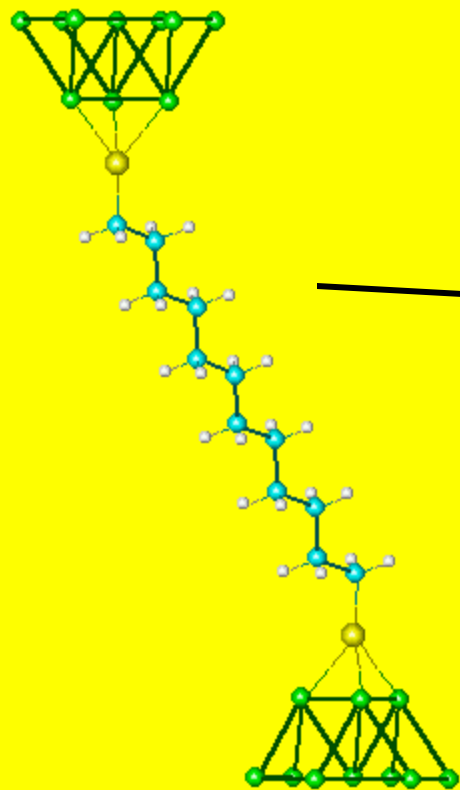
Simulation of molecular devices

New physically-based simulations are needed that accounts for the electronic structure of the molecule => **Approach:** Non equilibrium Green Function (based on a Density Functional Tight Binding method) calculation selfconsistently coupled to a 3D Poisson solver

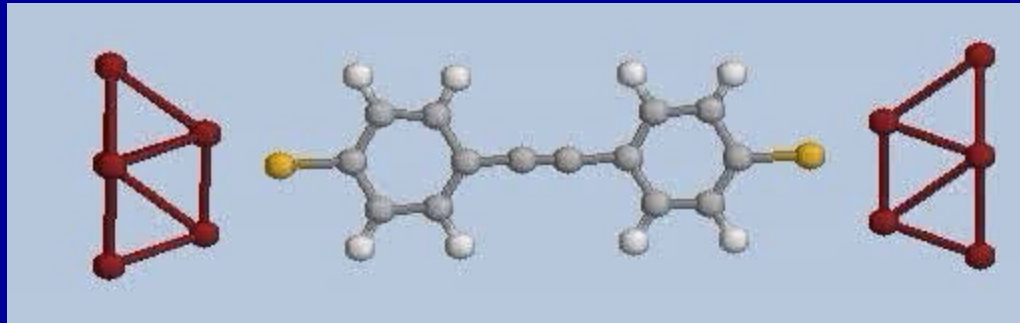


Tunneling currents in alkeno-thiols

Exp. M. Reed et al. (2003)



The role of molecular vibrations



$T = 300 \text{ K}$

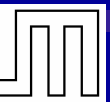
An organic molecule is a rather “floppy” entity

Electron-phonon interactions are a difficult issue

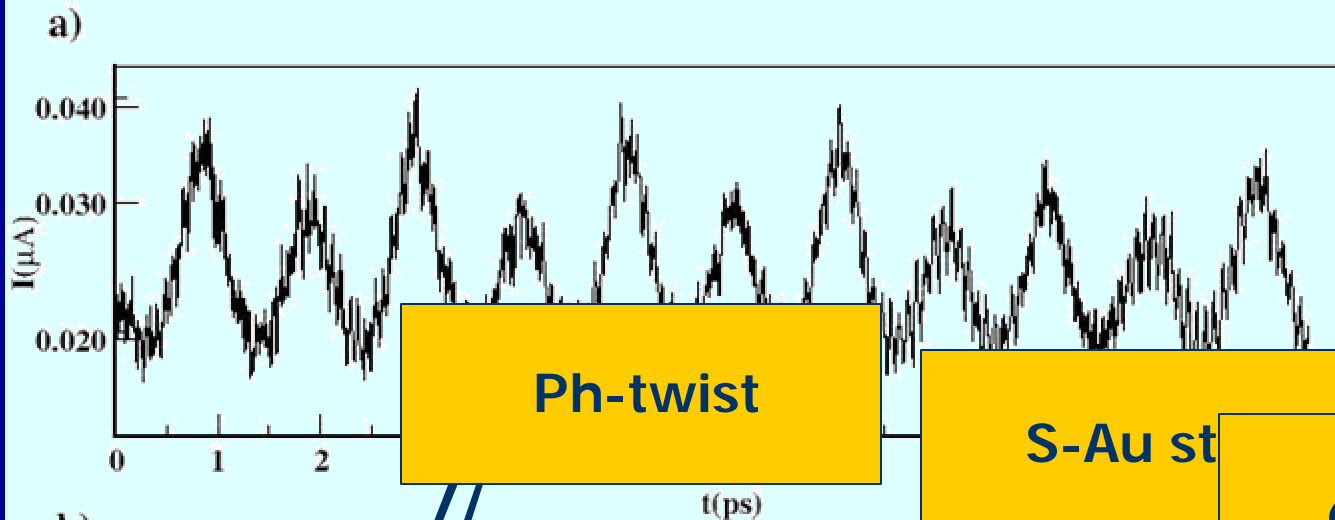
Can we still use the phonon approximation ?

Two approaches:

- Time-average of the current computed at every step of a MD simulation
- Ensemble average of over the lattice fluctuations.

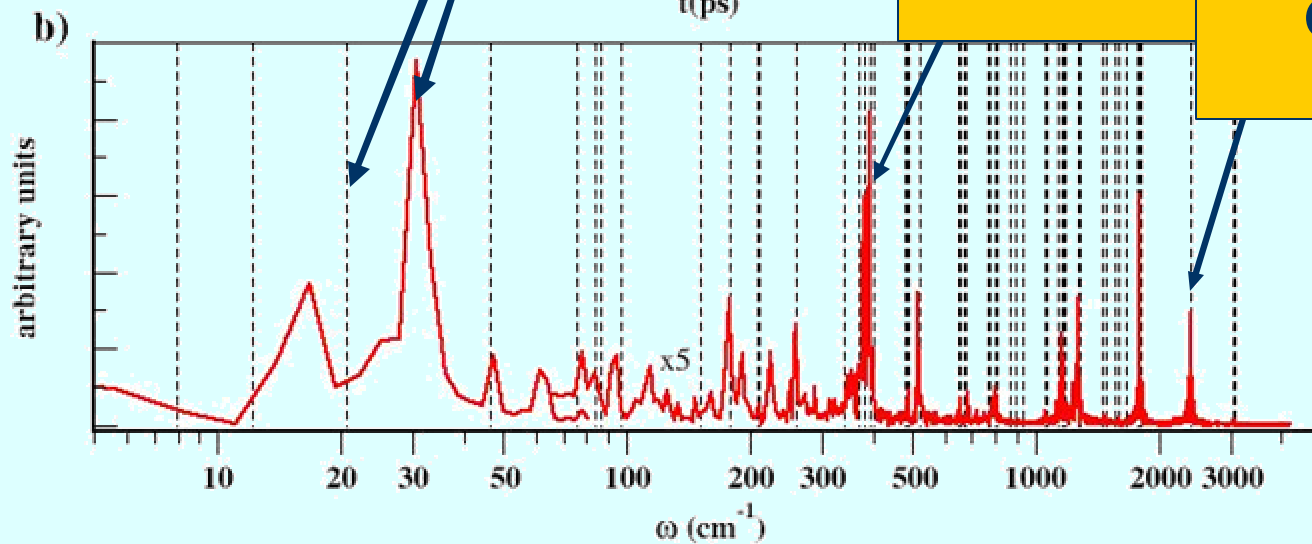


Frequency analysis



Mol. Dynamic

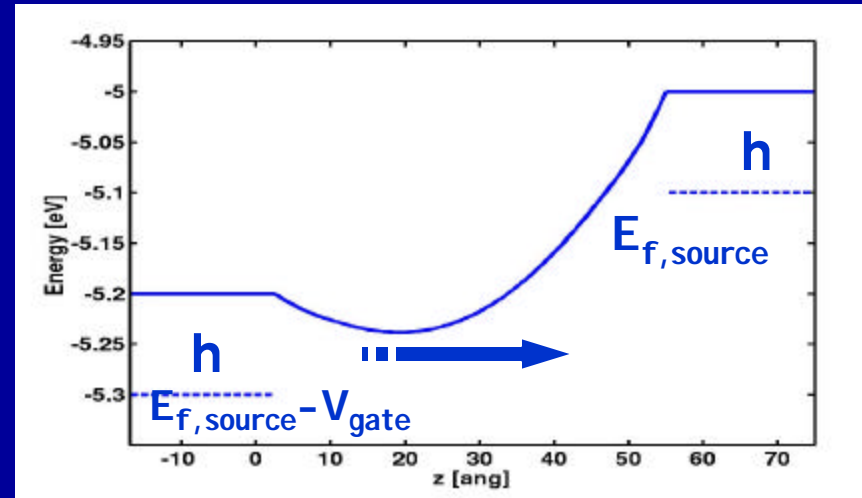
$T=10\text{ K}$



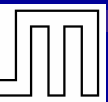
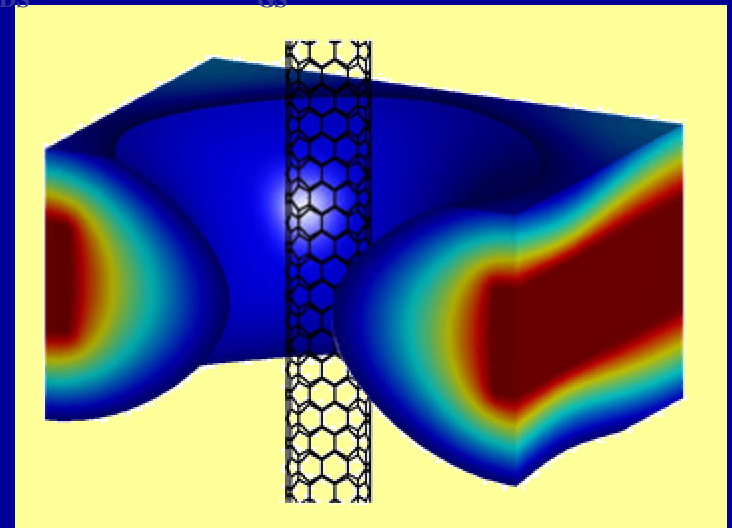
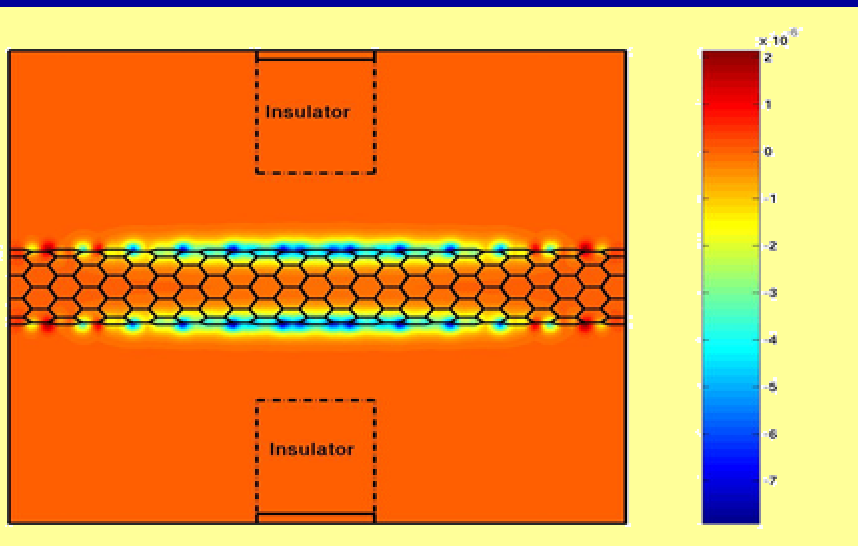
Fourier Trans

Simulation of CNT-Transistors

The CNTs behave like ideal p-type nanowires. The gate voltage modulates the transmission of the holes which are injected from the source contact.



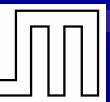
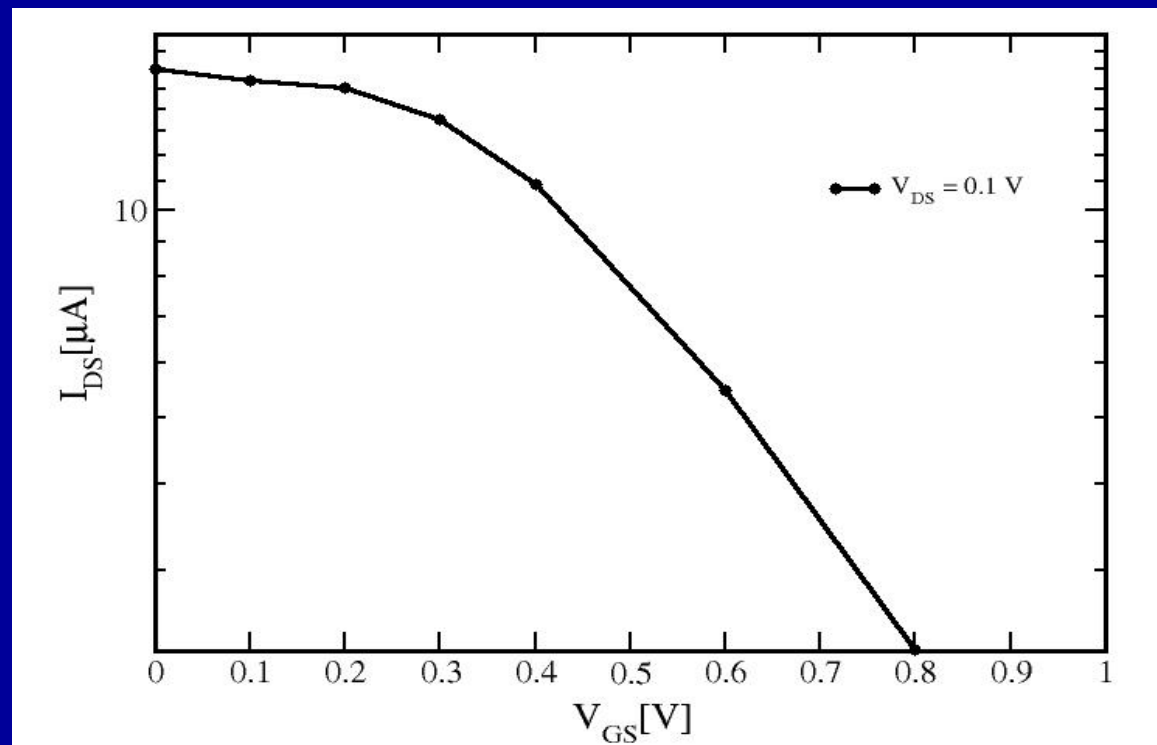
Das Energie-Profil des Valenz-Subbandes einer (10,0) CNT für $V_{DS} = 0.2$ V und $V_{GS} = 0.6$ V.



Current in CNT-Transistors

Landauer Formalism is used to calculate the current between source and drain

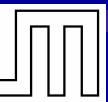
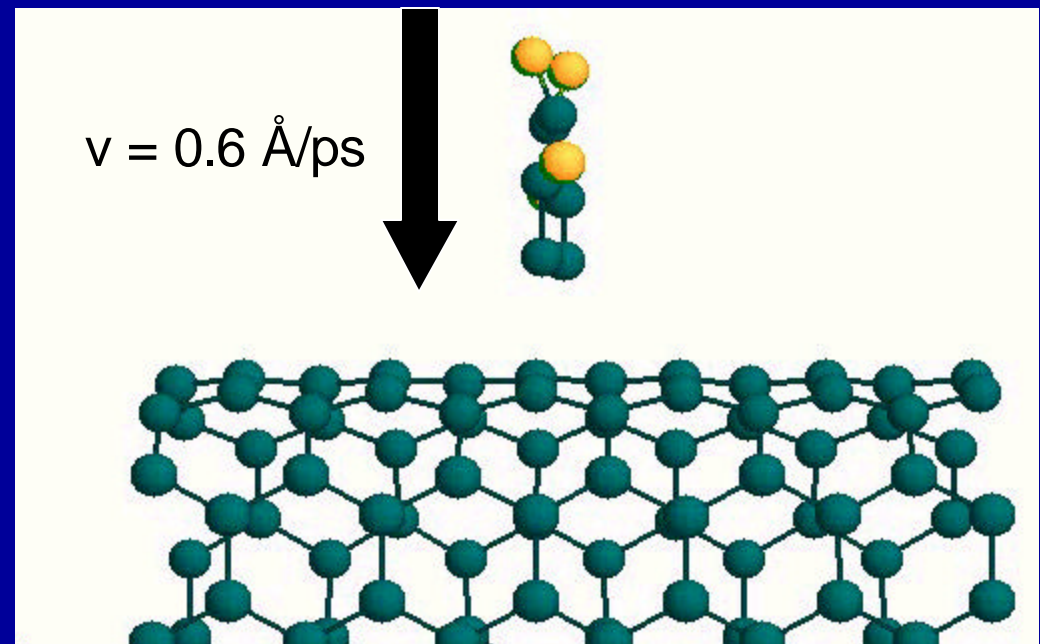
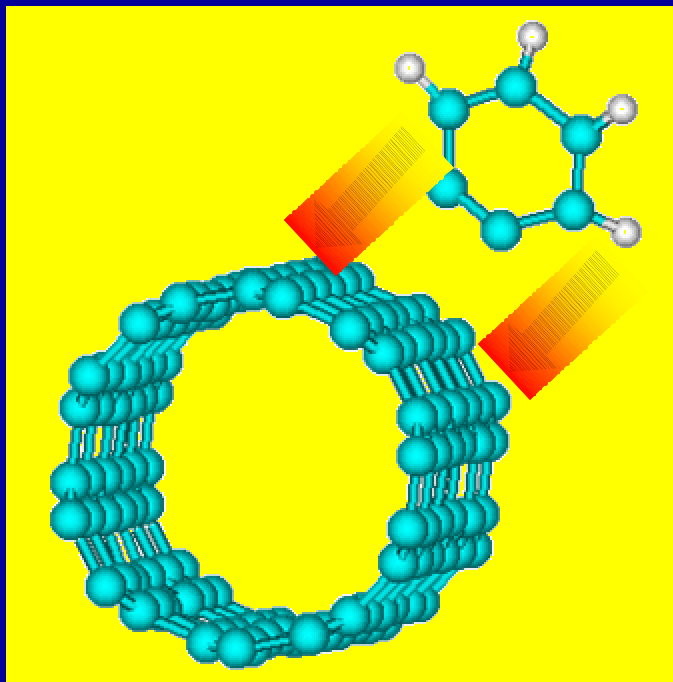
$$I = \frac{2e}{h} \int \bar{T}(E) [f_1^+(E) - f_2^-(E)] dE$$

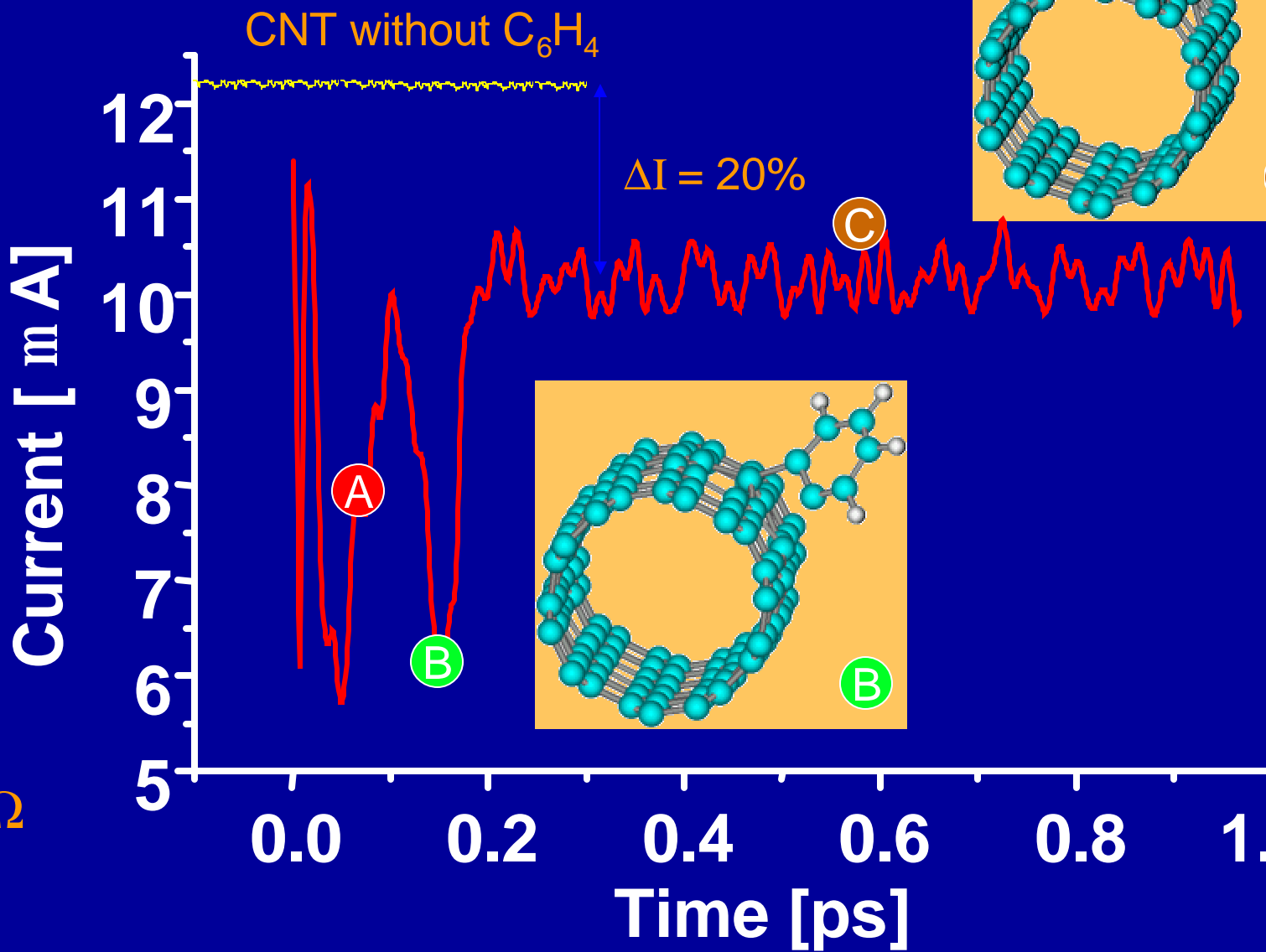
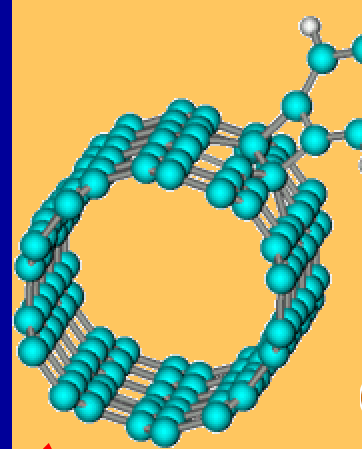
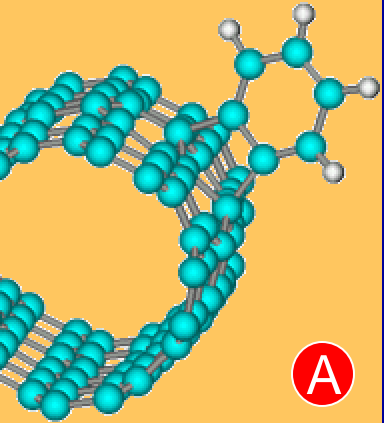


Transport in Carbon Nanotube Green + MD

We have performed Molecular Dynamics simulations of a reactive collision of a biased nanotube ($V=100\text{mV}$) and benzyne and we have calculated the current flowing in the nanotube at each MD step

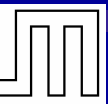
[Following the NASA MD simulations (J. Han, A. Globus, R. Jaffe, G. Deardorff)]



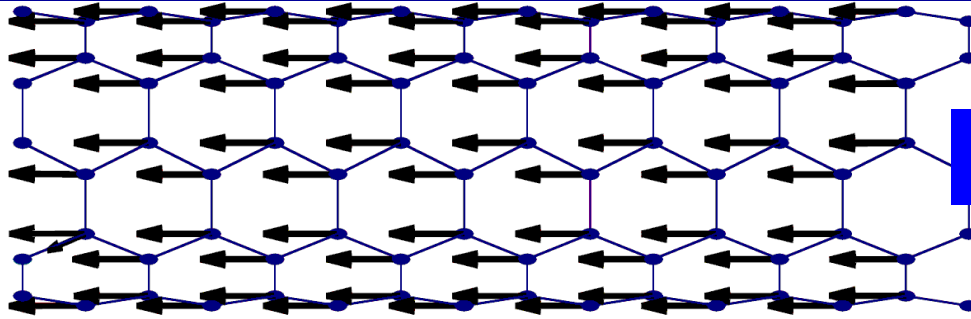


$$R_{CN-C_6H_4} = 10 K\Omega$$

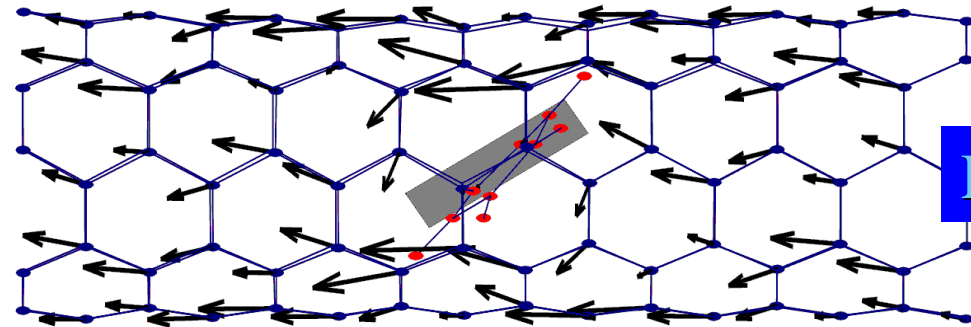
$$R_{CN} = 8 K\Omega$$



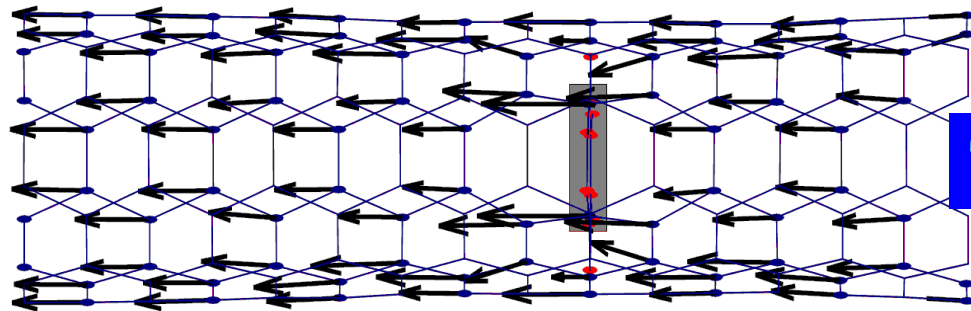
Current Flux in CNT (5,5)



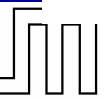
Unperturbed Nanotube



Longitudinal Molecule

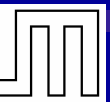


Transverse Molecule



Open problems

- ❖ the theoretical treatment of excited states in molecular system;
- ❖ the interaction with phonons and in general the dissipative part of transport;
- ❖ the description of the contact and the molecule-contact interface;
- ❖ the integration of models properly accounting for excited states, dissipation and contacts in a simulator
- ❖ the identification of optimal circuits and architectures for molecular devices
- ❖ the availability and reliability of experimental results

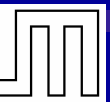


Architectures

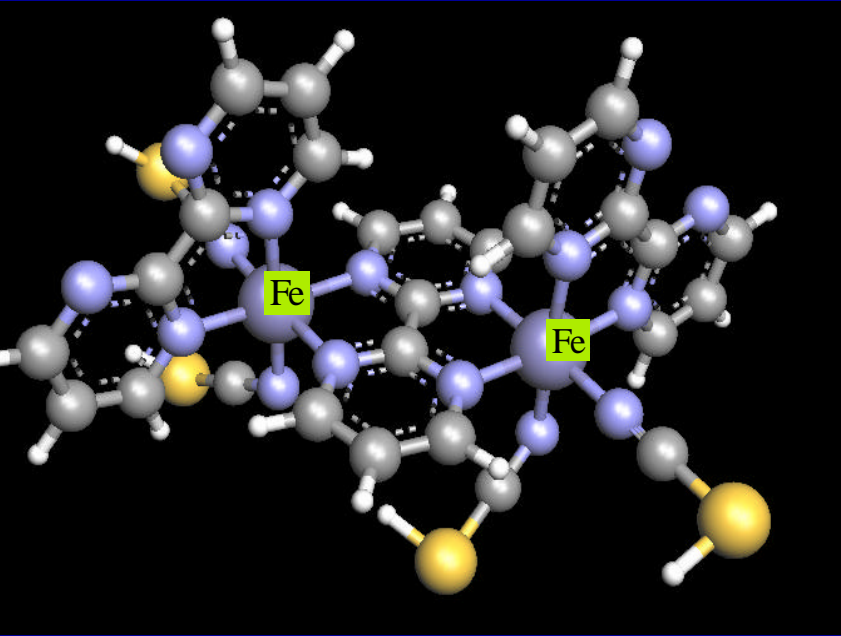
Single molecule devices require not only novel simulation tools, but also novel architectures that can exploit the interaction between molecules and the high parallelism intrinsic in the molecular systems

Some suggestions:

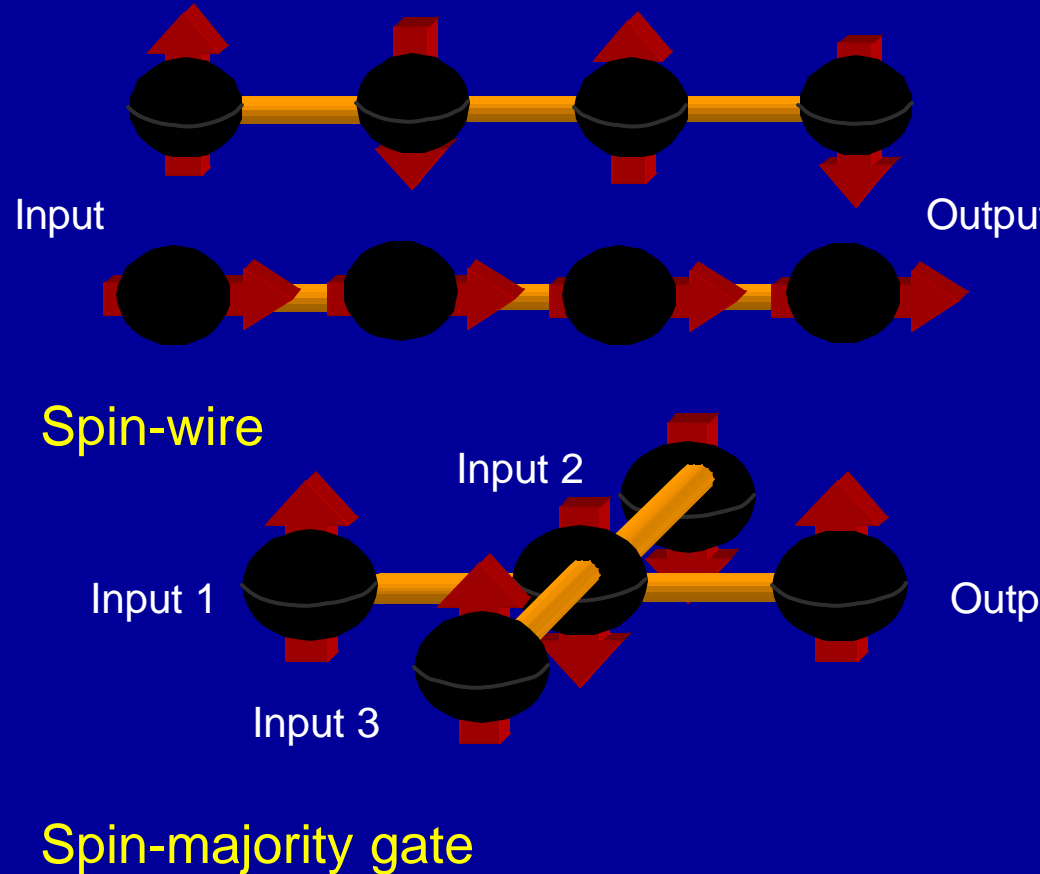
- ❖ Molecular Logic gates (J. Tour, MITRE Corp.)
- ❖ Crossbar memory (HP, STMicroelectronics)
- ❖ Quantum Cellular Automata (Notre Dame)



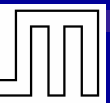
Field-Coupling by Molecular Magnets



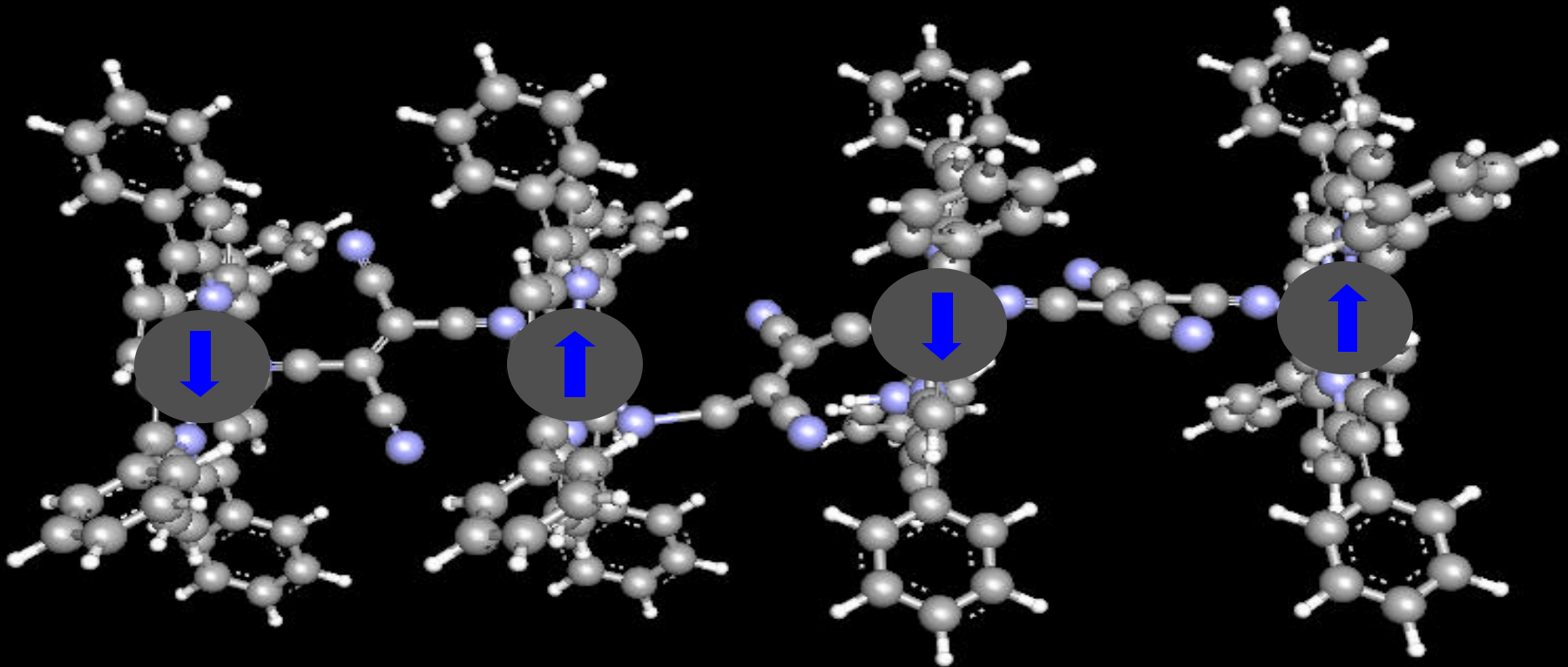
$[\text{Fe}(\text{bpym})(\text{NCS})]_2\text{bpym}$



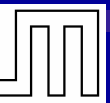
In molecular magnets, exchange interaction between metal ions is mediated by organic ligands. Certain arrangements could make QCA-like logic gates.



A 1D Molecular Magnet Chain



Manganese ions in a $[\text{Mn}^{\text{III}}(\text{Porphyrin})][\text{TCNE}]$ chain might play the role of lithographically fabricated nanoscale magnets. This would yield to a molecular-scale computing device.



Conclusions

- ❖ The simulation of molecular structures has progressed greatly in the last few years
- ❖ Future advances require an even stronger integration of quantum chemistry and transport codes
- ❖ As reliable experiments will become available, validation of the models will be possible;
- ❖ Circuit design will be performed with the same codes used to analyzed molecular nanostructures

